

Turbulent Flows
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Solution to Exercise 12.25

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a): According to the definition of the fluctuating velocity following a particle, we get

$$\mathbf{u}^*(t) \equiv \mathbf{U}^*(t) - \langle \mathbf{U}(\mathbf{X}^*[t], t) \rangle. \quad (1)$$

Supposing $\mathbf{X}^*(t)$ evolves by

$$\frac{d\mathbf{X}^*(t)}{dt} = \mathbf{U}^*(t), \quad (2)$$

and $\mathbf{U}^*(\mathbf{t})$ evolves by a diffusion process, from Eq. 1, we get

$$\begin{aligned} du_i^*(t) &= dU_i^*(t) - d\langle U_i(\mathbf{X}^*[t], t) \rangle \\ &= dU_i^*(t) - \left[\frac{\partial \langle U_i \rangle}{\partial t} + \left(\frac{\partial \langle U_i \rangle}{\partial x_j} \right)_{\mathbf{x}=\mathbf{X}^*} \frac{dX_j^*}{dt} \right] dt \\ &= dU_i^*(t) - \left(\frac{\partial \langle U_i \rangle}{\partial t} + \frac{dX_j^*}{dt} \frac{\partial \langle U_i \rangle}{\partial x_j} \right) dt \\ &= dU_i^*(t) - \left[\frac{\partial \langle U_i \rangle}{\partial t} + (u_j^*(t) + \langle U_j \rangle) \frac{\partial \langle U_i \rangle}{\partial x_j} \right] dt \\ &= dU_i^*(t) - u_j^* \frac{\partial \langle U_i \rangle}{\partial x_j} dt - \frac{\bar{D} \langle U_i \rangle}{\bar{D}t} dt. \end{aligned} \quad (3)$$

b): For homogeneous turbulence, and for $\mathbf{U}^*(t)$ evolving by (GLM)

$$dU_i^*(t) = -\frac{1}{\rho} \frac{\partial \langle p \rangle}{\partial x_i} dt + G_{ij}(U_j^*(t) - \langle U_j \rangle) dt + (C_0 \varepsilon)^{\frac{1}{2}} dW_i(t), \quad (4)$$

we get

$$du_i^*(t) = -\frac{1}{\rho} \frac{\partial \langle p \rangle}{\partial x_i} dt + G_{ij}(U_j^*(t) - \langle U_j \rangle) dt + (C_0 \varepsilon)^{\frac{1}{2}} dW_i(t)$$

$$\begin{aligned}
& -u_j^* \frac{\partial \langle U_i \rangle}{\partial x_j} dt - \frac{\bar{D} \langle U_i \rangle}{\bar{D}t} dt \\
= & -\frac{1}{\rho} \frac{\partial \langle p \rangle}{\partial x_i} dt + G_{ij} (U_j^*(t) - \langle U_j \rangle) dt + (C_0 \varepsilon)^{\frac{1}{2}} dW_i(t) \\
& -u_j^* \frac{\partial \langle U_i \rangle}{\partial x_j} dt + \frac{1}{\rho} \frac{\partial \langle p \rangle}{\partial x_i} dt - \nu \nabla^2 \langle U_i \rangle dt \\
= & -u_j^* \frac{\partial \langle U_i \rangle}{\partial x_j} dt + G_{ij} u_j^* dt + (C_0 \varepsilon)^{\frac{1}{2}} dW_i(t), \tag{5}
\end{aligned}$$

where the second step follows from the mean momentum equation and the last step follows from homogeneity.

c): According to the Fokker-Plank equation, the evolution equation for $g_L^*(\mathbf{v}, \mathbf{x}; t | \mathbf{Y})$ can be derived from Eq 2 and Eq. 5. The result is

$$\begin{aligned}
\frac{\partial g_L^*}{\partial t} + \frac{\partial}{\partial x_i} [g_L^* (\langle U_i \rangle + v_i)] &= \frac{\partial}{\partial v_i} \left[g_L^* v_j \frac{\partial \langle U_i \rangle}{\partial x_j} \right] - \frac{\partial}{\partial v_i} [G_{ij} v_j g_L^*] \\
&+ \frac{1}{2} (C_0 \varepsilon) \frac{\partial^2 g_L^*}{\partial v_i \partial v_i}. \tag{6}
\end{aligned}$$

Using $\nabla \cdot \langle \mathbf{U} \rangle = 0$, from Eq. 6, we get

$$\begin{aligned}
\frac{\partial g_L^*}{\partial t} + (\langle U_i \rangle + v_i) \frac{\partial g_L^*}{\partial x_i} &= v_j \frac{\partial \langle U_i \rangle}{\partial x_j} \frac{\partial g_L^*}{\partial v_i} - \frac{\partial}{\partial v_i} [G_{ij} v_j g_L^*] \\
&+ \frac{1}{2} (C_0 \varepsilon) \frac{\partial^2 g_L^*}{\partial v_i \partial v_i}. \tag{7}
\end{aligned}$$

Integrating Eq. 7 for g_L^* over all \mathbf{Y} , we get the evolution equation for $g^*(\mathbf{x}; t)$

$$\begin{aligned}
\frac{\partial g^*}{\partial t} + (\langle U_i \rangle + v_i) \frac{\partial g^*}{\partial x_i} &= v_j \frac{\partial \langle U_i \rangle}{\partial x_j} \frac{\partial g^*}{\partial v_i} - \frac{\partial}{\partial v_i} [G_{ij} v_j g^*] \\
&+ \frac{1}{2} (C_0 \varepsilon) \frac{\partial^2 g^*}{\partial v_i \partial v_i}. \tag{8}
\end{aligned}$$

For homogeneous turbulence, the second term on the left-hand side of Eq. 8 vanishes, so Eq. 8 turns into

$$\frac{\partial g^*}{\partial t} - v_j \frac{\partial \langle U_i \rangle}{\partial x_j} \frac{\partial g^*}{\partial v_i} = -\frac{\partial}{\partial v_i} [G_{ij} v_j g^*] + \frac{1}{2} (C_0 \varepsilon) \frac{\partial^2 g^*}{\partial v_i \partial v_i}, \tag{9}$$

and this is identical to Eq.(12.27).

d): According to Eq. 9, we get

$$\frac{\partial g^*}{\partial t} = v_j \frac{\partial \langle U_i \rangle}{\partial x_j} \frac{\partial g^*}{\partial v_i} - \frac{\partial}{\partial v_i} [G_{ij} v_j g^*] + \frac{1}{2} (C_0 \varepsilon) \frac{\partial^2 g^*}{\partial v_i \partial v_i}. \quad (10)$$

Multiplying Eq. 10 by $v_i v_j$ and integrating over the fluctuating velocity space, we get

$$\begin{aligned} \int v_i v_j \frac{\partial g^*}{\partial t} d\mathbf{v} &= \int v_i v_j v_k \frac{\partial \langle U_\ell \rangle}{\partial x_k} \frac{\partial g^*}{\partial v_\ell} d\mathbf{v} - \int v_i v_j \frac{\partial}{\partial v_k} (G_{k\ell} v_\ell g^*) d\mathbf{v} \\ &+ \int v_i v_j \frac{1}{2} C_0 \varepsilon \frac{\partial^2 g^*}{\partial v_k \partial v_k} d\mathbf{v}. \end{aligned} \quad (11)$$

From Eq. 11, it is easy to get

$$\begin{aligned} \frac{\partial \langle u_i^* u_j^* \rangle}{\partial t} &= -\langle u_j^* u_k^* \rangle \frac{\partial \langle U_i \rangle}{\partial x_k} - \langle u_i^* u_k^* \rangle \frac{\partial \langle U_j \rangle}{\partial x_k} + G_{j\ell} \langle u_\ell^* u_i^* \rangle \\ &+ G_{i\ell} \langle u_\ell^* u_j^* \rangle + C_0 \varepsilon \delta_{ij}, \end{aligned} \quad (12)$$

and this is consistent with Eq.(12.57). In particular, the first term on the right-hand side of Eq. 5 leads to the production.

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